Fast Methods for Simulation of Biomolecule Electrostatics

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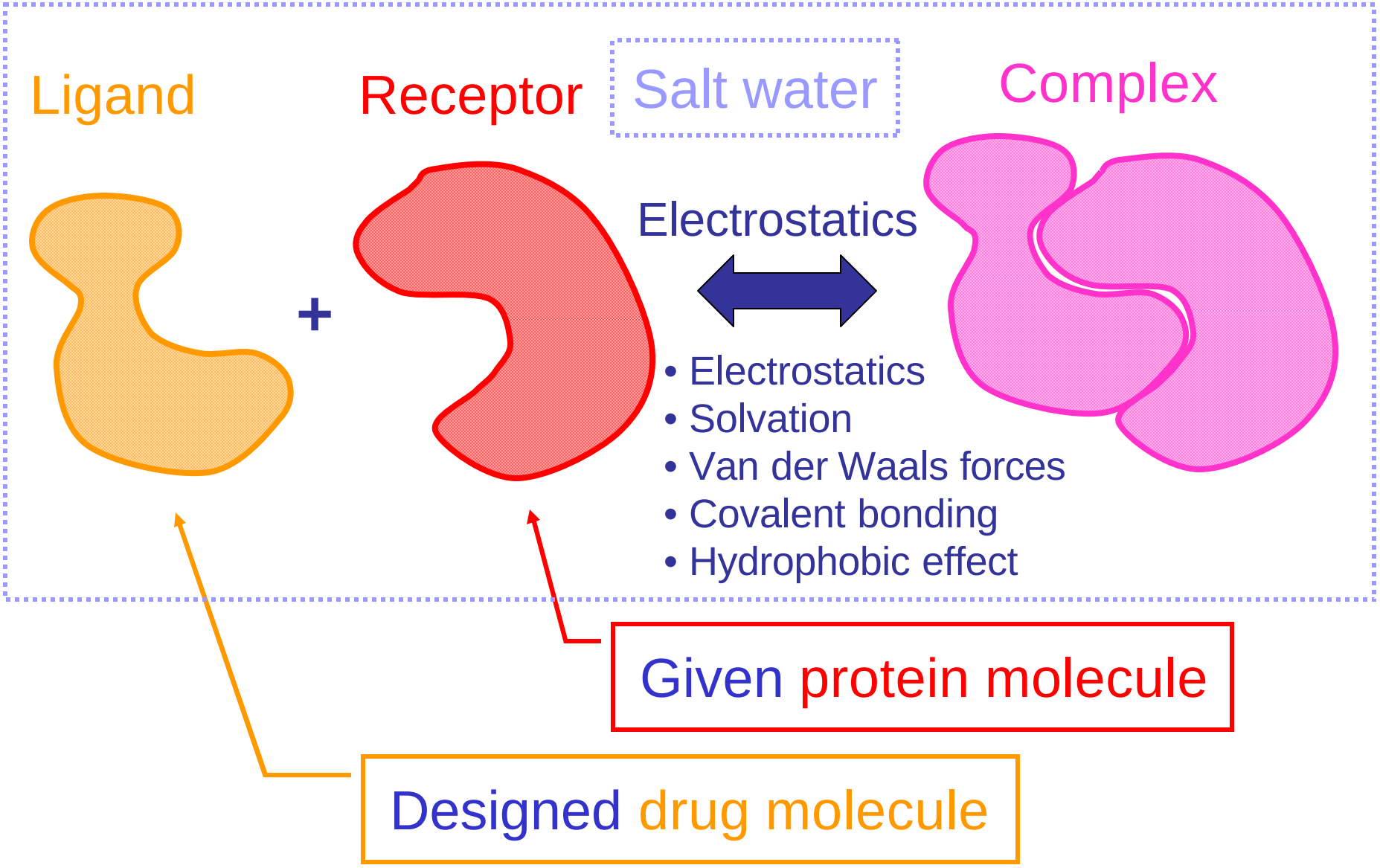
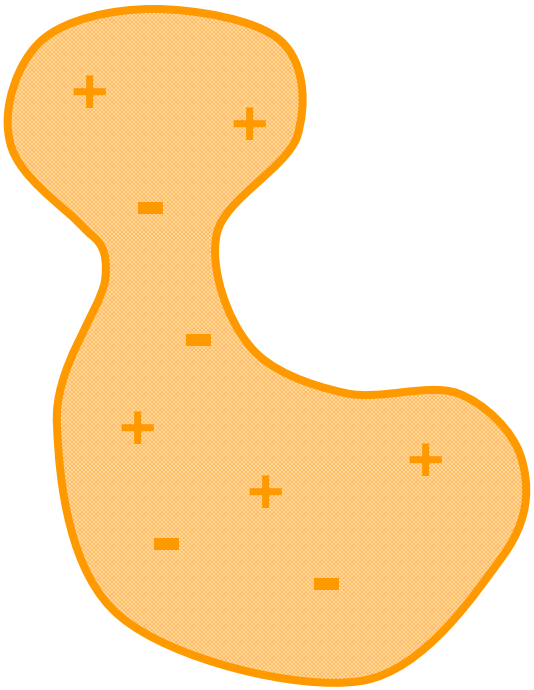
Outline

- Problem statement
- Finite-difference approach and problems
- Integral equation method and advantages
- Fast solver implementation
- Computational results
- Conclusion and Future work

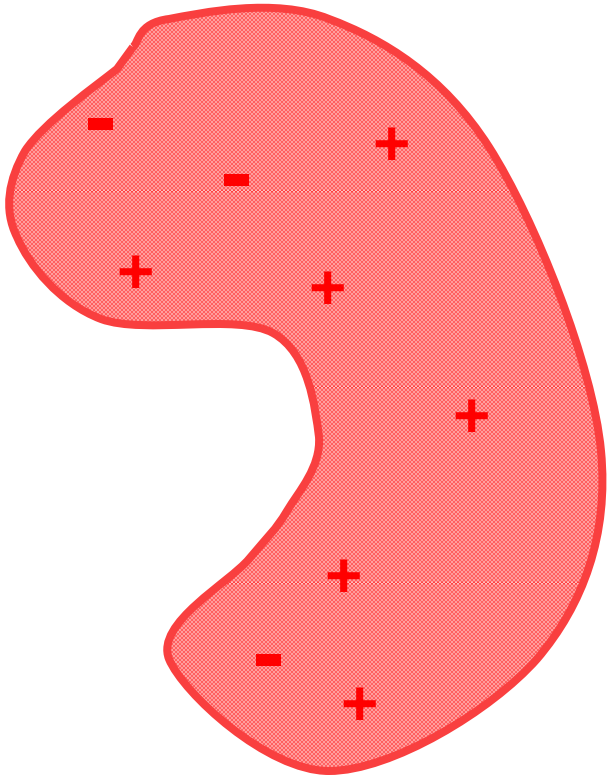
The Problem of Drug Design



Electrostatics View



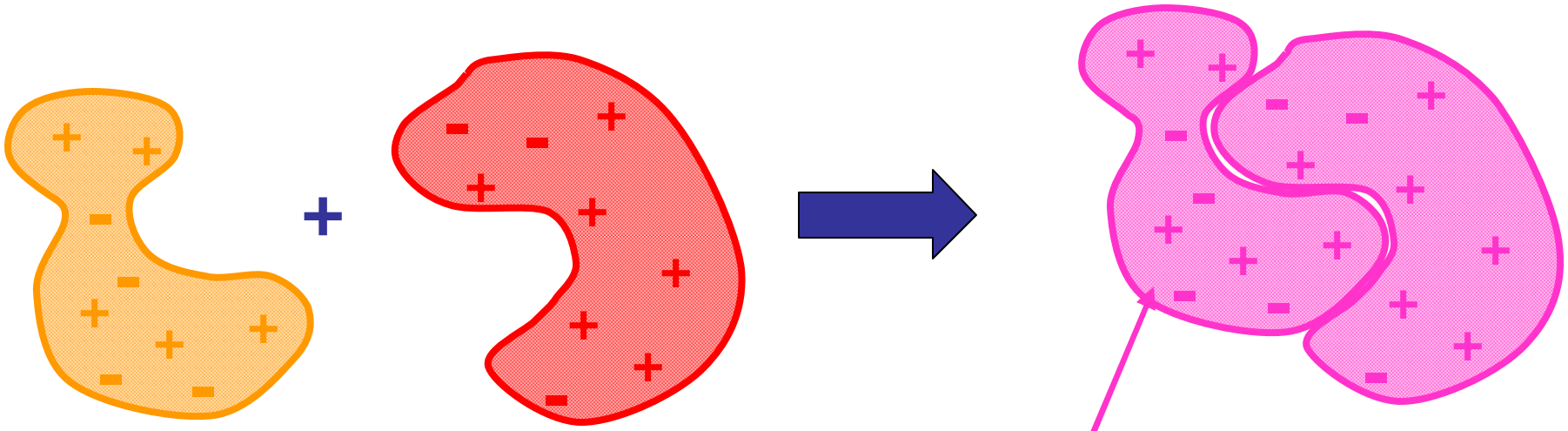
Ligand
(drug molecule)



Receptor
(protein molecule)

Minimize Electrostatic Binding Energy

$$E_{binding} = \left| E^{ligand}_{desolvation} \right| + \left| E^{receptor}_{desolvation} \right| - \left| E_{interaction} \right|$$

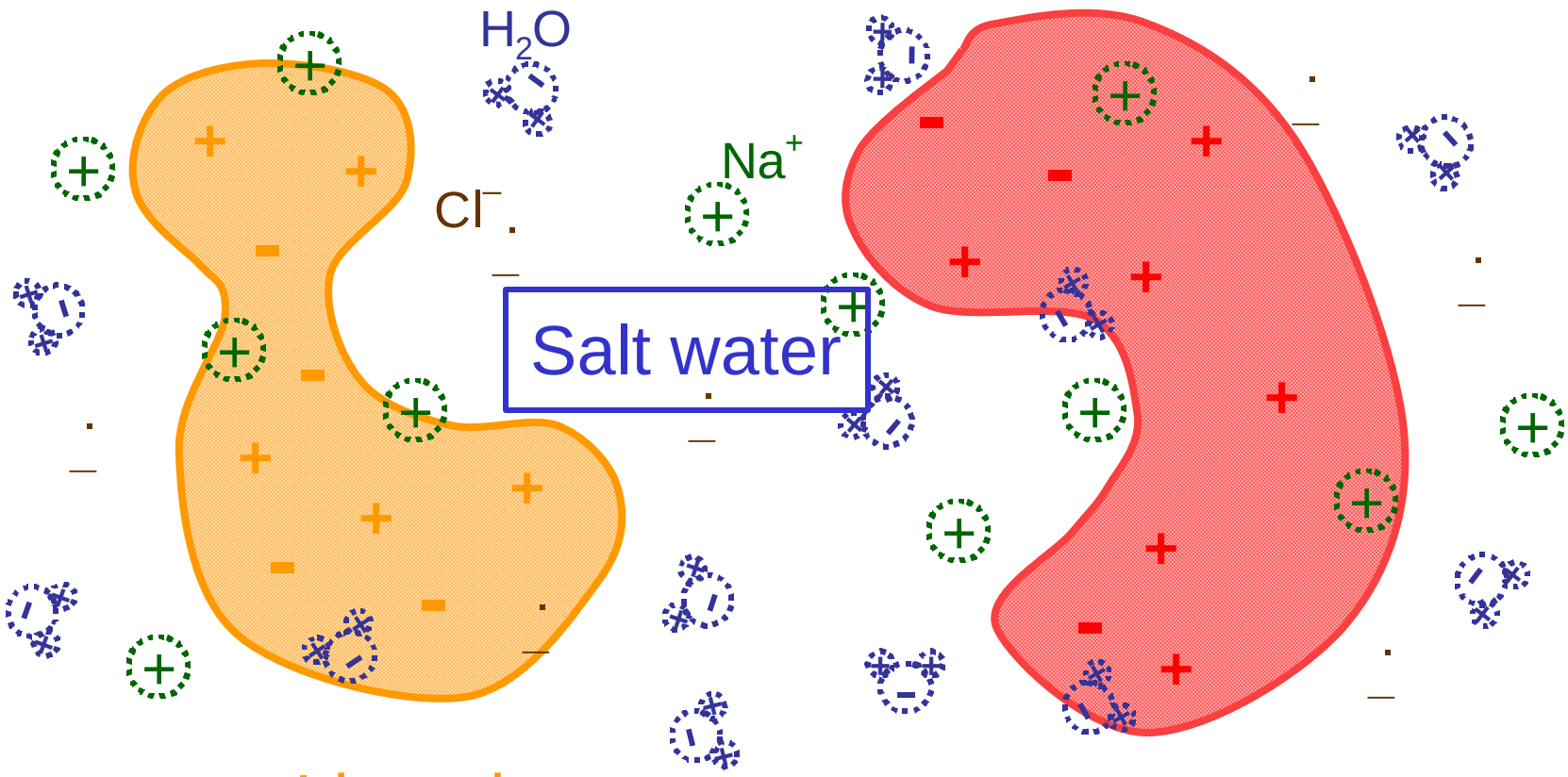


Higher energy

Lower energy

Determine the charge distribution in the ligand so that it is “Energetically Optimized” to bind

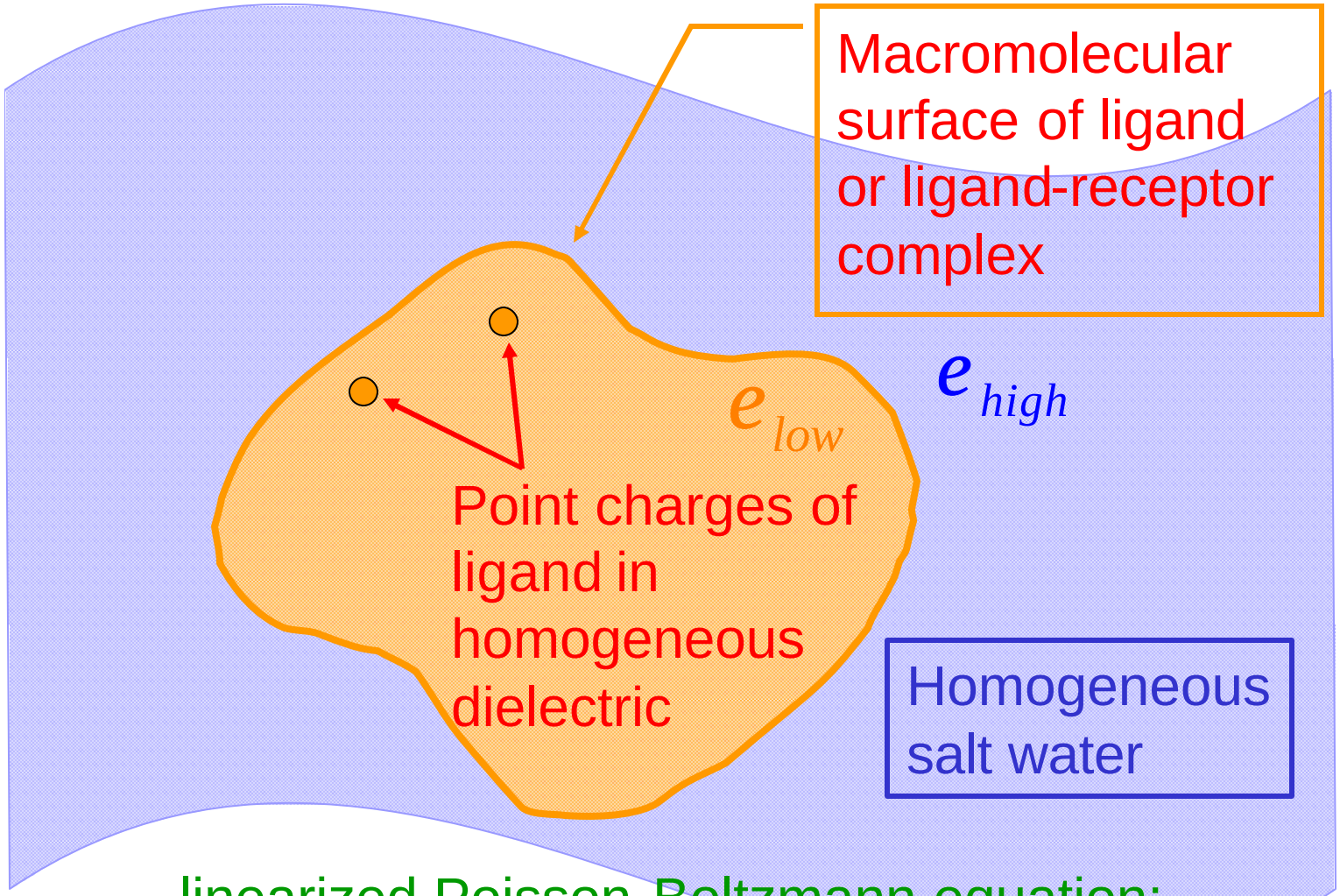
An Electrostatic Analysis Problem



Ligand
(drug molecule)

Receptor
(protein molecule)

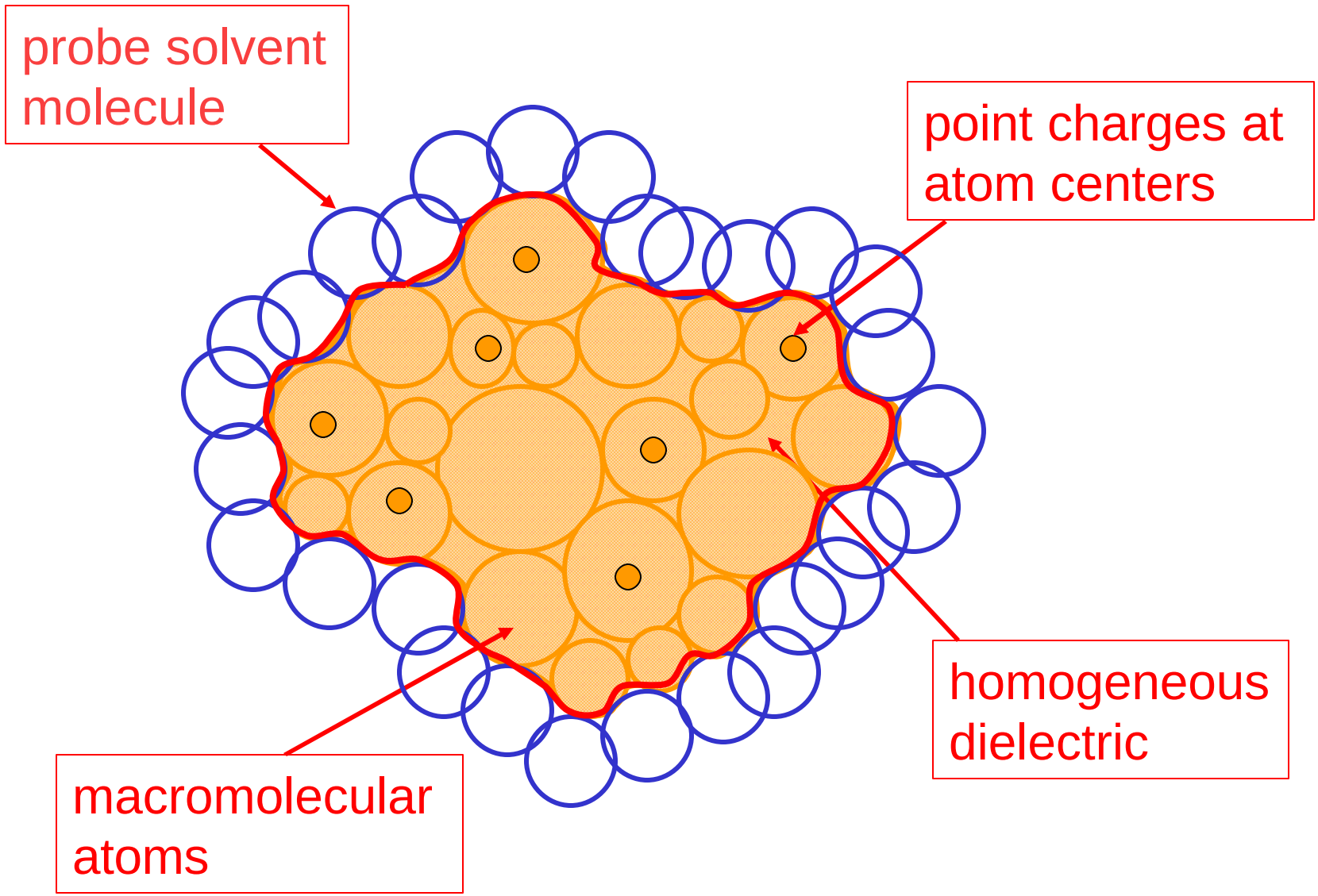
A Simplified Physical Description



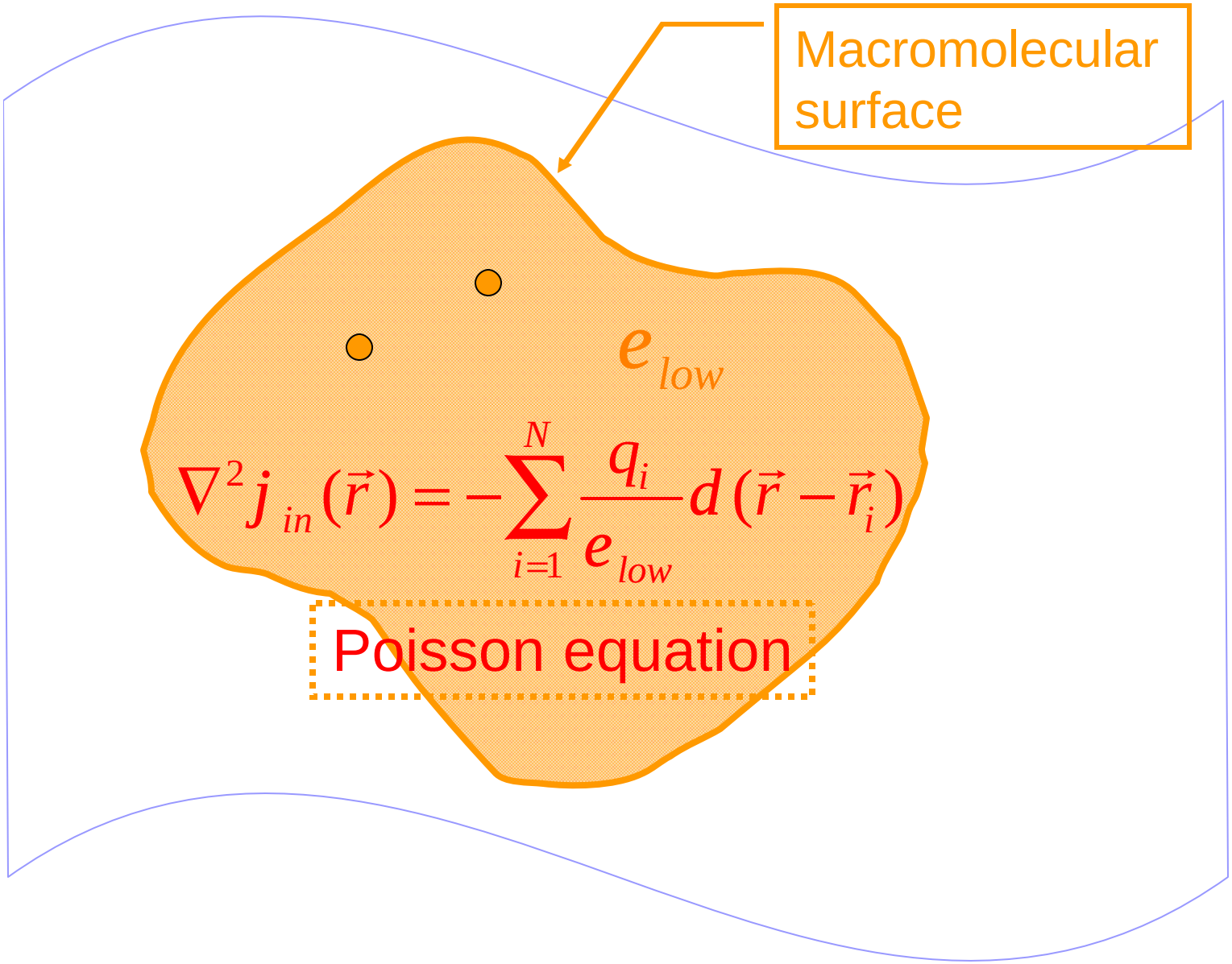
linearized Poisson-Boltzmann equation:

$$\nabla \bullet \left(e(\vec{r}) \nabla j(\vec{r}) \right) - e(\vec{r}) k^2(\vec{r}) j(\vec{r}) + r(\vec{r}) = 0$$

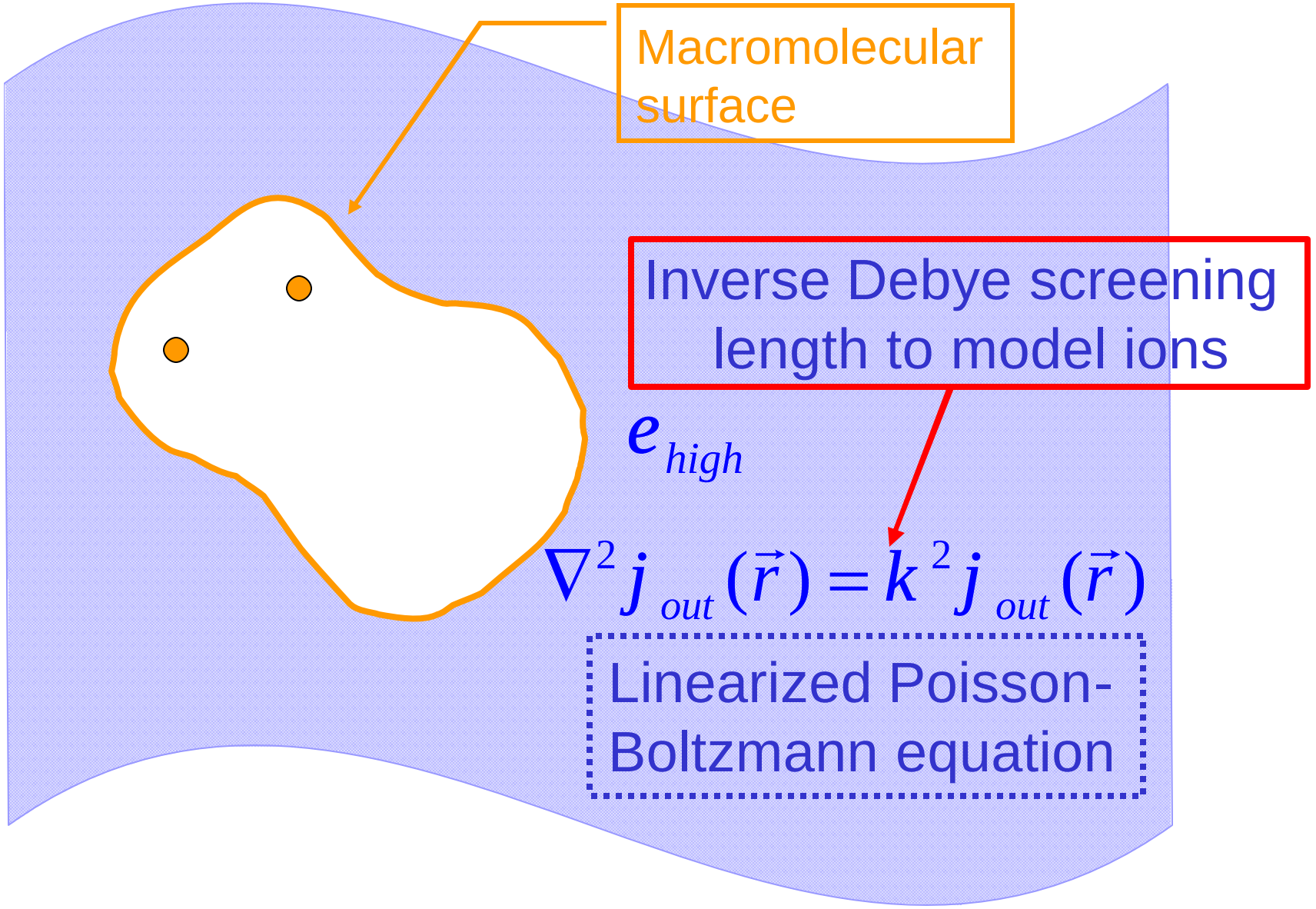
Molecular Surface Representation



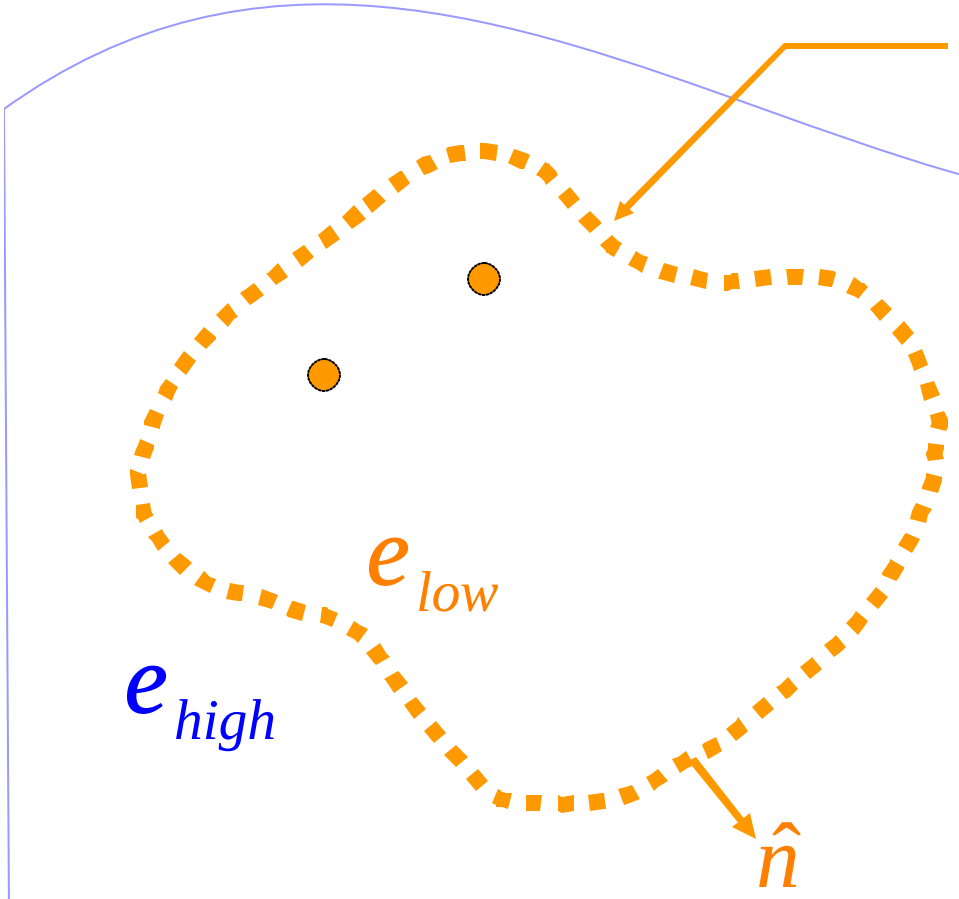
Simplified Mathematical Model: Inside Macromolecule



Simplified Mathematical Model: Salt Water Outside



Interface Condition



Boundary Conditions:

$$j_{in}(\vec{r}_o) = j_{out}(\vec{r}_o)$$

$$\frac{\partial j_{in}}{\partial n}(\vec{r}_o) = e_r \frac{\partial j_{out}}{\partial n}(\vec{r}_o)$$

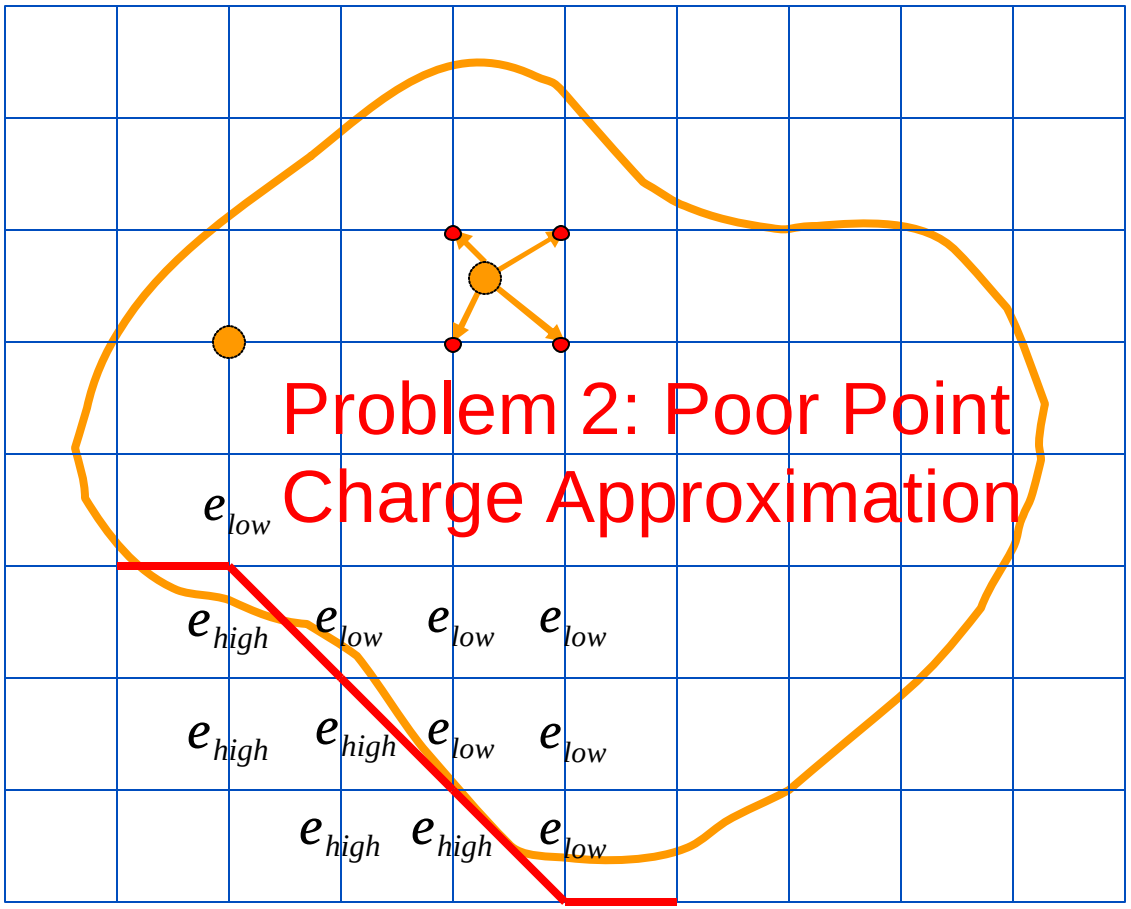
$$e_r = \frac{e_{high}}{e_{low}}$$

Why Use this Simplified Model?

- Atomistic Level Simulation is too expensive
  - Salt ions and water molecules treated individually
- Continuum Model Matches Well with Experimental Data

Standard Finite-Difference Method

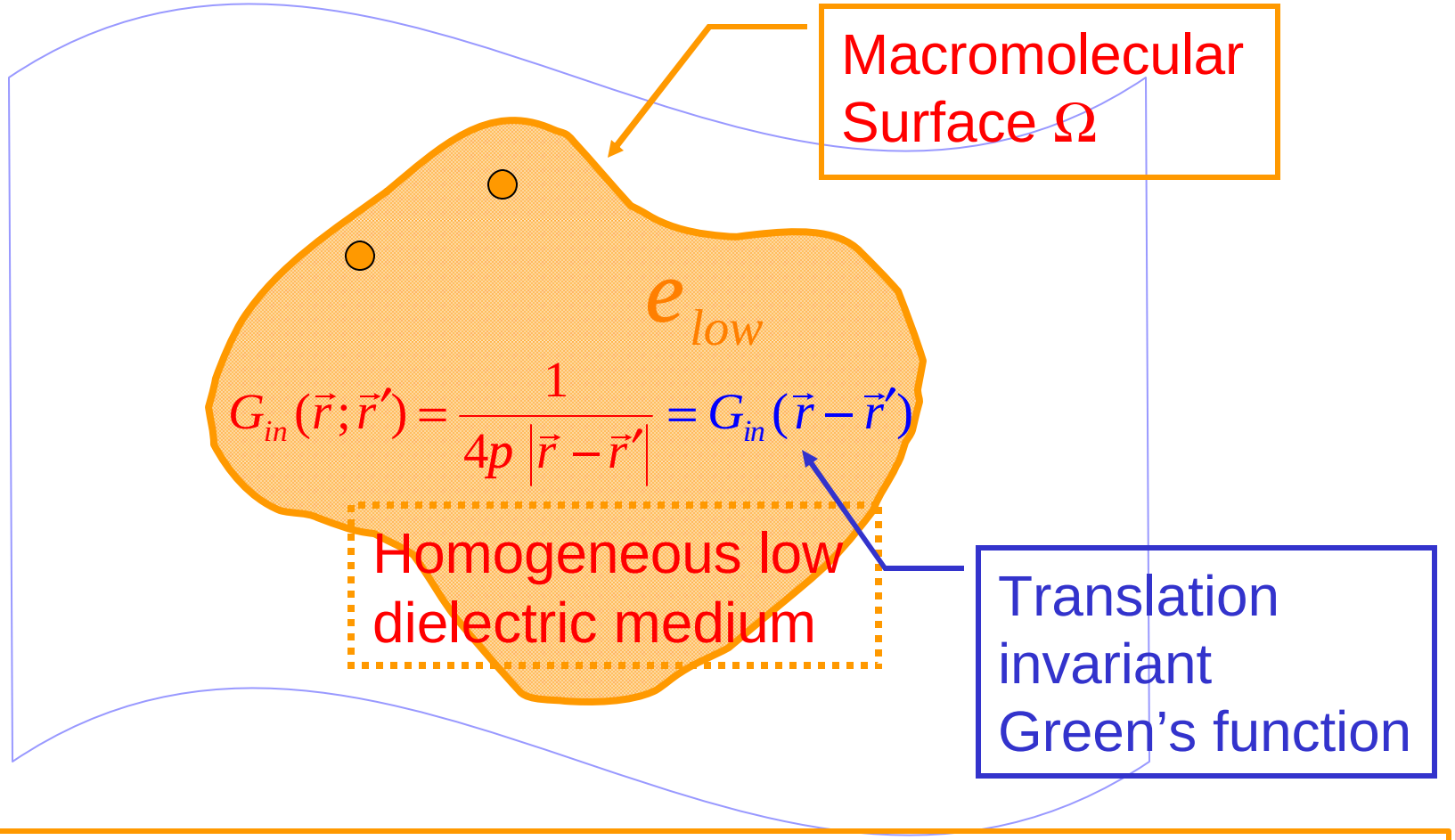
set up boundary conditions and
solve for grid potentials



Problem 3: Inexact Boundary Conditions

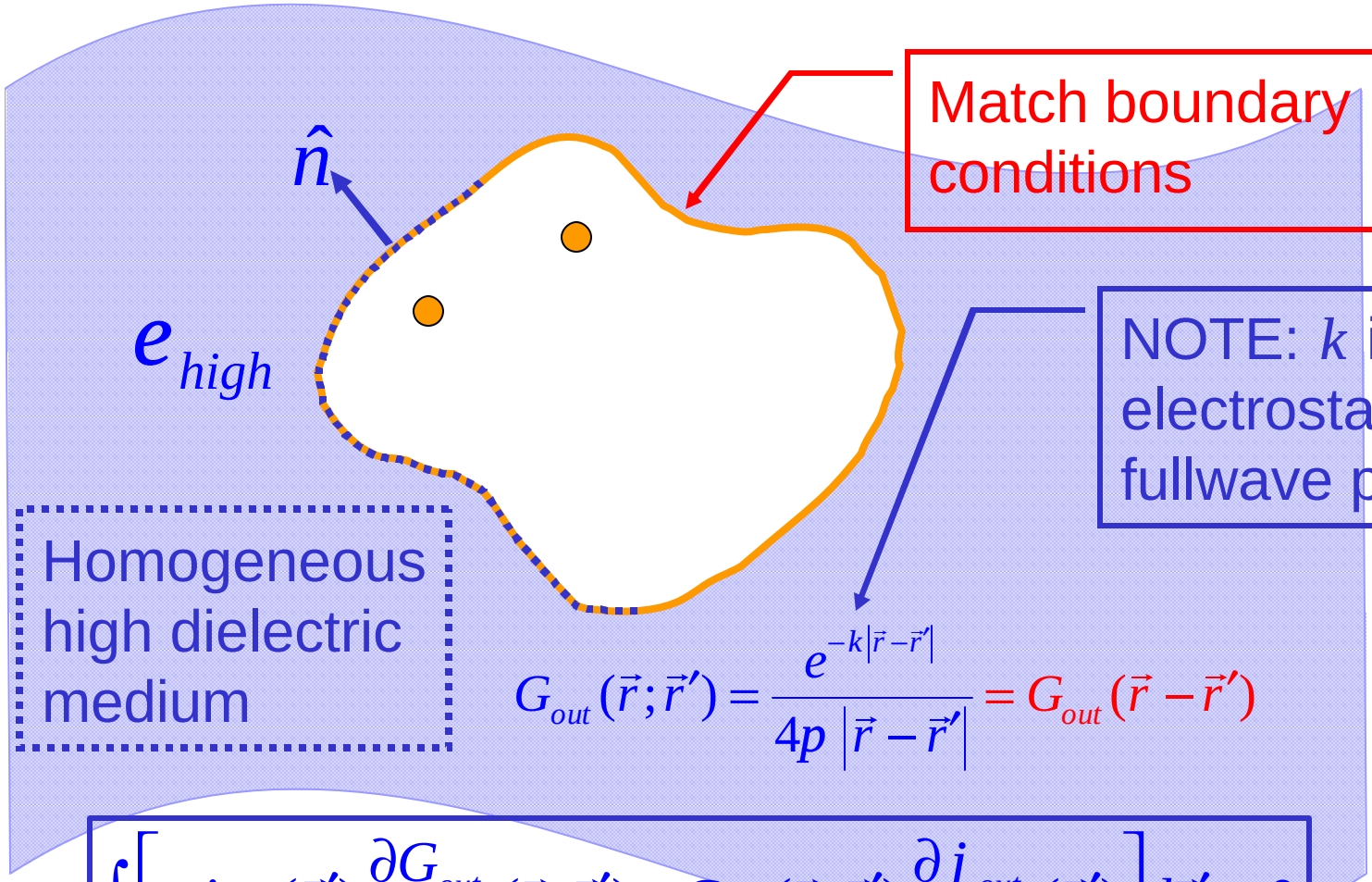
Problem 1: Inaccurate Molecular Surface

Integral equation: Interior Problem



$$\int_{\Omega}\left[j_{in}(\vec{r}')\frac{\partial G_{in}}{\partial n}(\vec{r};\vec{r}') - G_{in}(\vec{r};\vec{r}')\frac{\partial j_{in}}{\partial n}(\vec{r}') \right] d\vec{r}' = \sum_{k=1}^{N} \frac{q_k}{e_{low}} G_{in}(\vec{r};\vec{r}_k)$$

Integral equation: Exterior Problem



Match boundary conditions

NOTE: k is real, electrostatic not fullwave problem

Homogeneous high dielectric medium

$$G_{out}(\vec{r};\vec{r}') = \frac{e^{-k|\vec{r}-\vec{r}'|}}{4p\,|\vec{r}-\vec{r}'|} = G_{out}(\vec{r}-\vec{r}')$$

$$\int_{\Omega}\left[-j_{out}(\vec{r}')\frac{\partial G_{out}}{\partial n}(\vec{r};\vec{r}') + G_{out}(\vec{r};\vec{r}')\frac{\partial j_{out}}{\partial n}(\vec{r}')\right]d\vec{r}' = 0$$

$j_{out}(\vec{r}') \rightarrow j_{in}(\vec{r}')$

$\frac{\partial j_{out}}{\partial n}(\vec{r}') \rightarrow \frac{1}{e_r}\frac{\partial j_{in}}{\partial n}(\vec{r}')$

Advantages For Integral Equation Formulation

- Directly discretize surfaces
- Point charges treated exactly
- Handles infinite exterior

$$\int_{\Omega}\left[j_{in}(\vec{r}')\frac{\partial G_{in}}{\partial n}(\vec{r};\vec{r}') - G_{in}(\vec{r};\vec{r}')\frac{\partial j_{in}}{\partial n}(\vec{r}') \right] d\vec{r}' = \sum_{k=1}^{N} \frac{q_k}{e_{low}} G_{in}(\vec{r};\vec{r}_k)$$

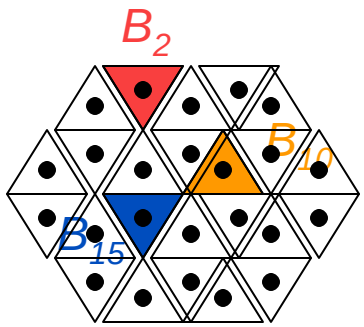
$$\int_{\Omega}\left[-j_{in}(\vec{r}')\frac{\partial G_{out}}{\partial n}(\vec{r};\vec{r}') + G_{out}(\vec{r};\vec{r}')\frac{1}{e_r}\frac{\partial j_{in}}{\partial n}(\vec{r}') \right] d\vec{r}' = 0$$

Standard piecewise constant collocation discretization method

$$j_{in}(\vec{r}) \approx \sum_j a_j B_j(\vec{r})$$

$$\frac{\partial j_{in}}{\partial n}(\vec{r}) \approx \sum_j b_j B_j(\vec{r})$$

$$\vec{r} \in \Omega$$



- Piecewise constant basis functions
- Collocation points at panel centroids

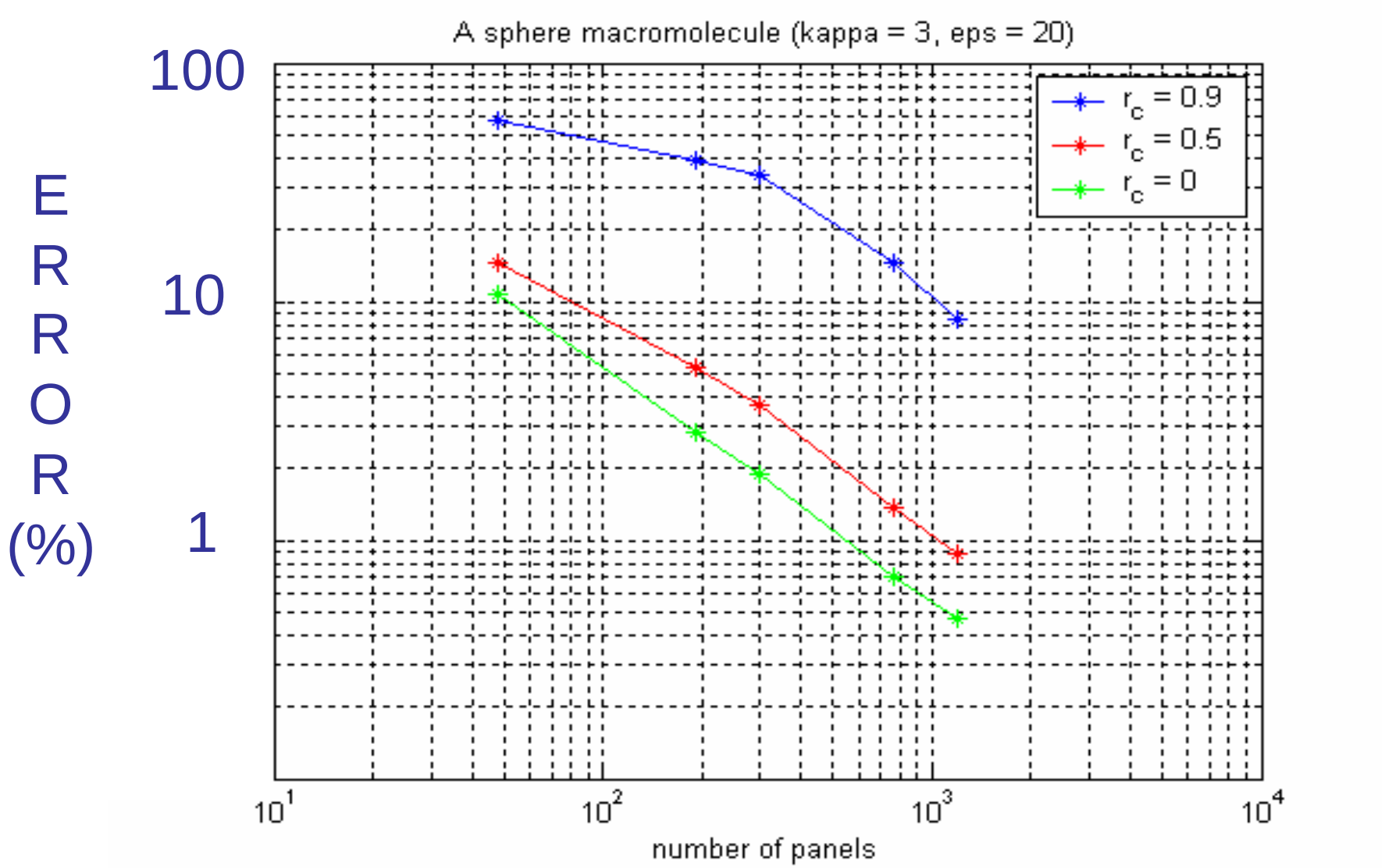
Matrix Equation

$$\begin{bmatrix} D^{in} & S^{in} \\ D^{out} & S^{out} \end{bmatrix} \begin{bmatrix} a_j \\ b_j \end{bmatrix} = \begin{bmatrix} \sum_{k=1}^{N} \frac{q_k}{4p\left|\vec{r}_i - \vec{r}_k\right|} \\ 0 \end{bmatrix}$$

$$D_{ij}^{in} = \int_{panel_j} \frac{\partial}{\partial n'} \left(\frac{1}{4p\left|\vec{r}_i - \vec{r}'\right|} \right) d\vec{r}' \qquad S_{ij}^{in} = -\int_{panel_j} \frac{1}{4p\left|\vec{r}_i - \vec{r}'\right|} d\vec{r}'$$

$$D_{ij}^{out} = -\int_{panel_j} \frac{\partial}{\partial n'} \left(\frac{e^{-k\left|\vec{r}_i - \vec{r}'\right|}}{4p\left|\vec{r}_i - \vec{r}'\right|} \right) d\vec{r}' \qquad S_{ij}^{out} = \frac{1}{e_r} \int_{panel_j} \frac{e^{-k\left|\vec{r}_i - \vec{r}'\right|}}{4p\left|\vec{r}_i - \vec{r}'\right|} d\vec{r}'$$

A sphere molecule: comparison with analytical result



Iterative solver

$$j(\vec{r}) \equiv \int K(\vec{r};\vec{r}')s(\vec{r}')d\vec{r}'$$

Discretization

$$[K_{i,j}][s_j] = [j_i]$$

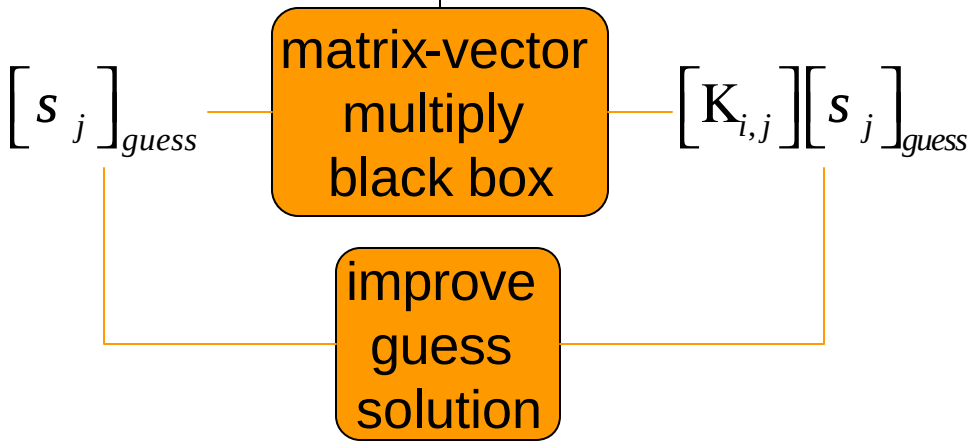
Gaussian Elimination

$$[s_j] = [K_{i,j}]^{-1}[j_i]$$

```
for i = 1:n-1
    for j = i+1:n
        K_{j,i} = K_{j,i} / K_{i,i}
        for k = i+1:n
            K_{j,k} = K_{j,k} - K_{j,i} / K_{i,k}
        end
    end
end
```

$O(N^3)$

Iterative Solver



$O(N^2)$ per iteration

Use Fast Integral Equation Solver

$O(N \log N)$ Matrix-vector multiply

- Multiple Green's functions
- Translation Invariant kernel

$$G_{in}(\vec{r};\vec{r}') = \frac{1}{4p\left|\vec{r}-\vec{r}'\right|}$$

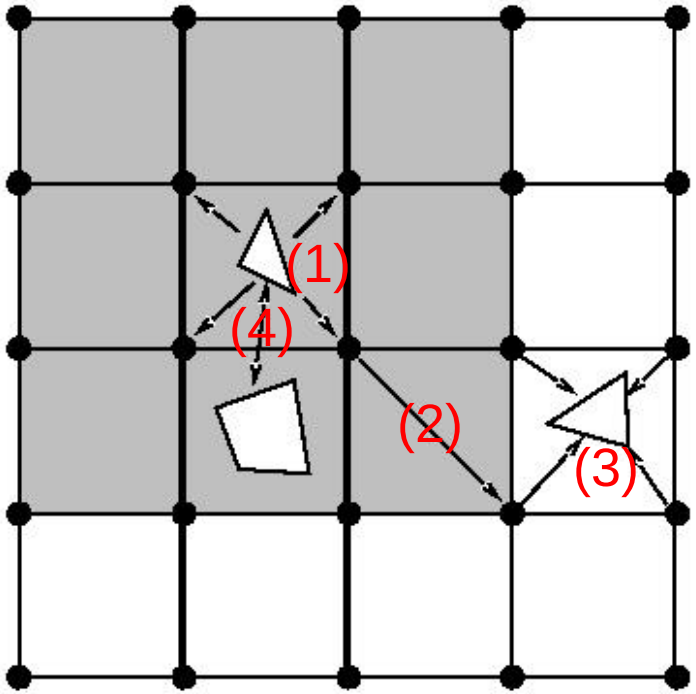
$$G_{out}(\vec{r};\vec{r}') = \frac{e^{-k\left|\vec{r}-\vec{r}'\right|}}{4p\left|\vec{r}-\vec{r}'\right|}$$

Pre-corrected FFT algorithm

$$\left[s_j \right]_{guess} \quad \text{matrix-vector multiply black box} \quad \left[K_{i,j} \right] \left[s_j \right]_{guess}$$

charge distribution — pre-corrected FFT — potential due to space-invariant kernel

1) projection of panel charges onto grid charges
2) grid potentials due to grid charges are computed by FFT
3) potentials on panel centroids are interpolated from grid potentials
4) direct interaction and correction among near neighbors



Picture courtesy of J. Phillips

Preconditioner on Two Examples

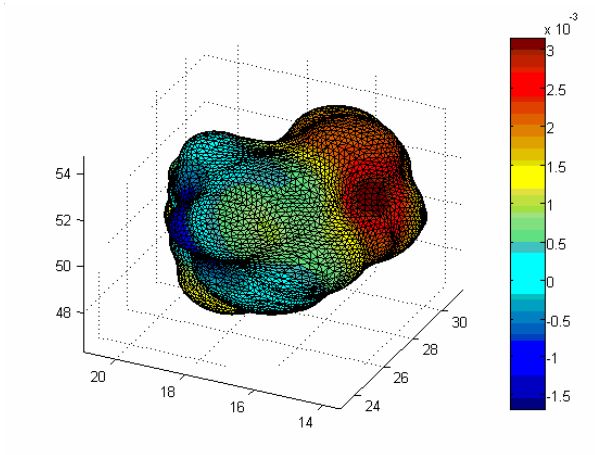
Need to find a good preconditioner

$$[P] \approx [\mathrm{K}_{i,j}]$$

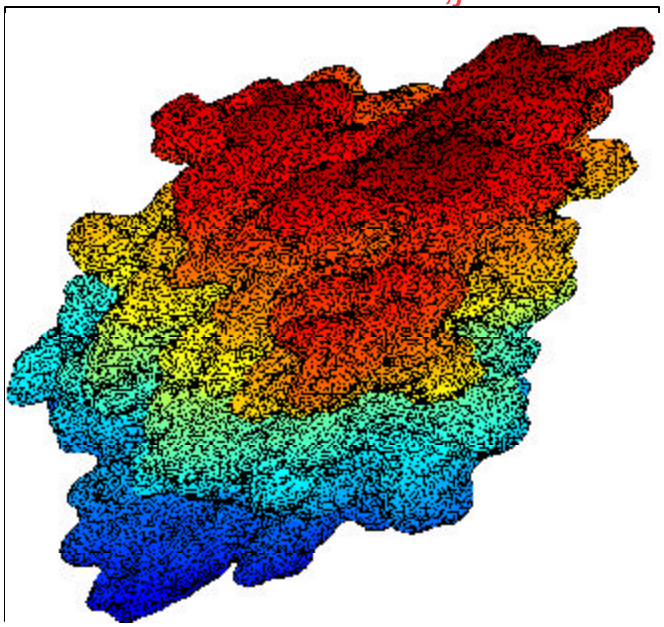
And solve

$$\underbrace{[P]^{-1}[\mathrm{K}_{i,j}]}[s_j] = [P]^{-1}[j_i]$$

hopefully better conditioned than $[K_{i,j}]$

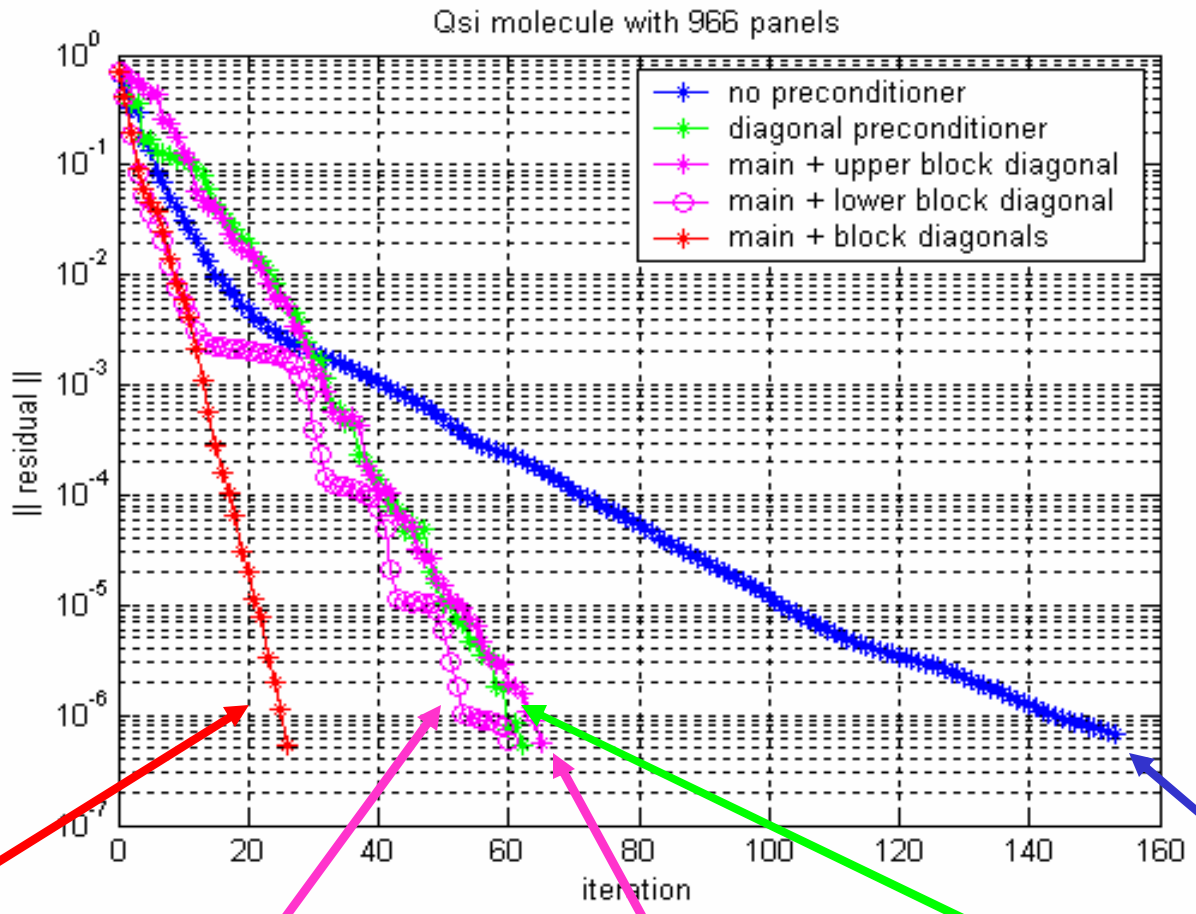


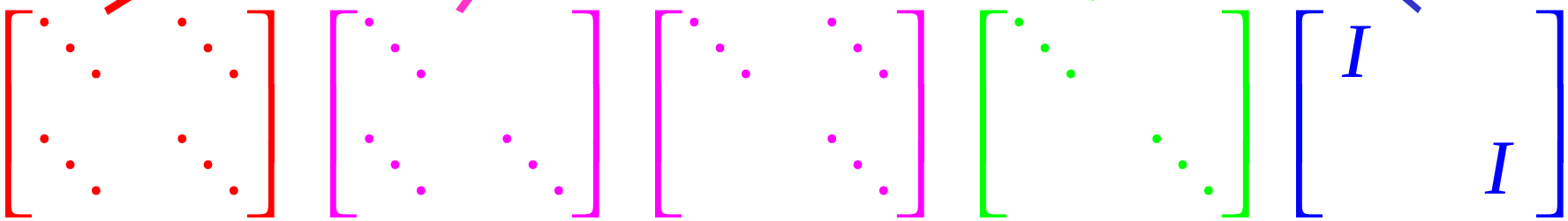
Qsi molecule



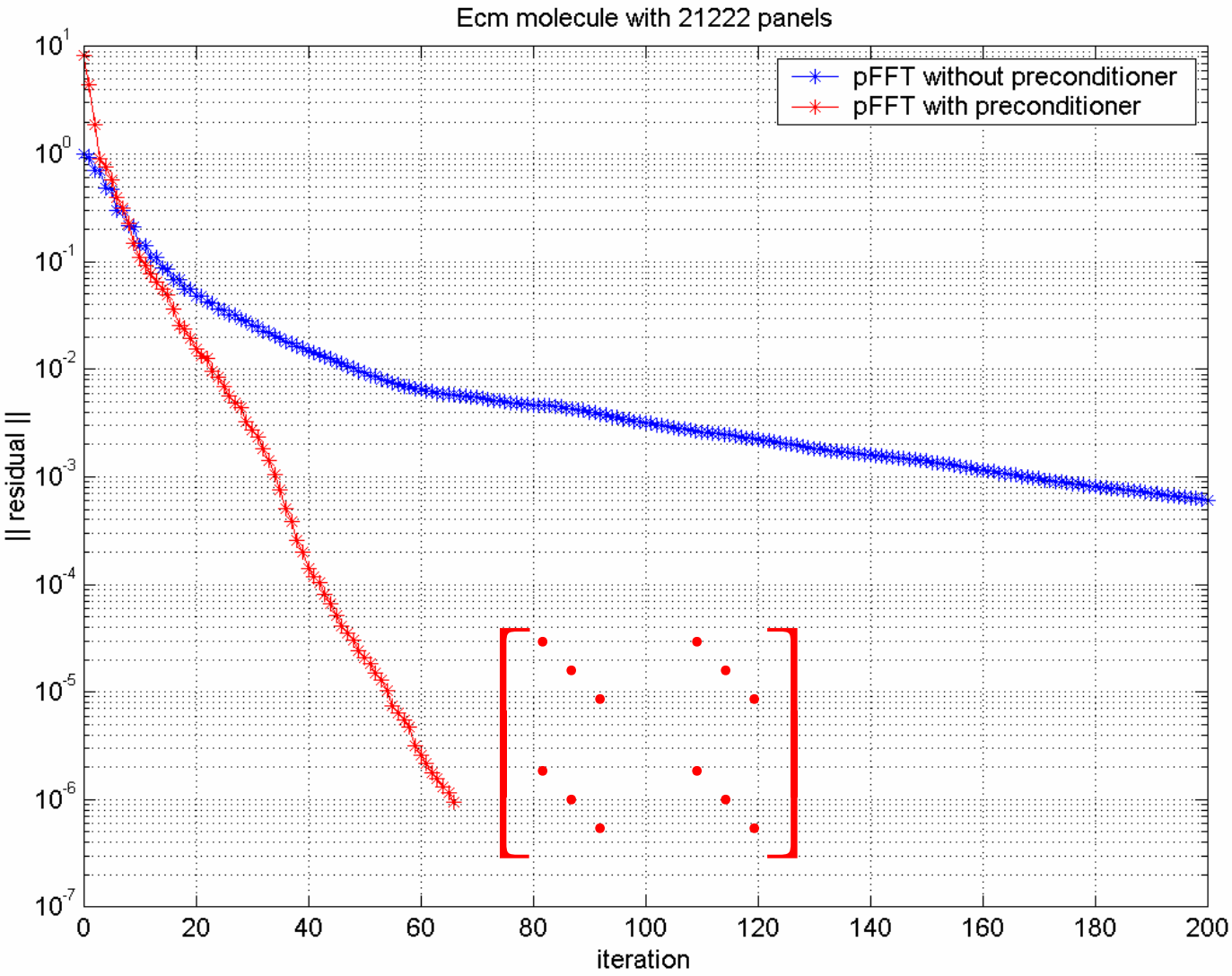
Ecm protein

Preconditioner result: Qsi molecule





Preconditioner result: Ecm protein

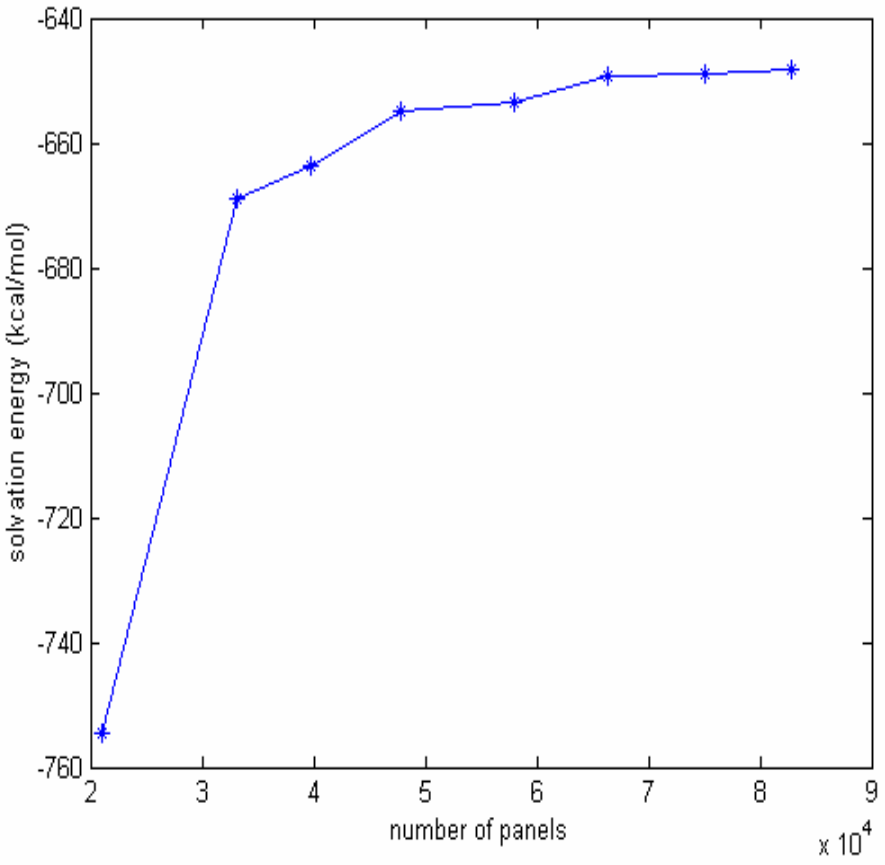


Accuracy comparison with DelPhi

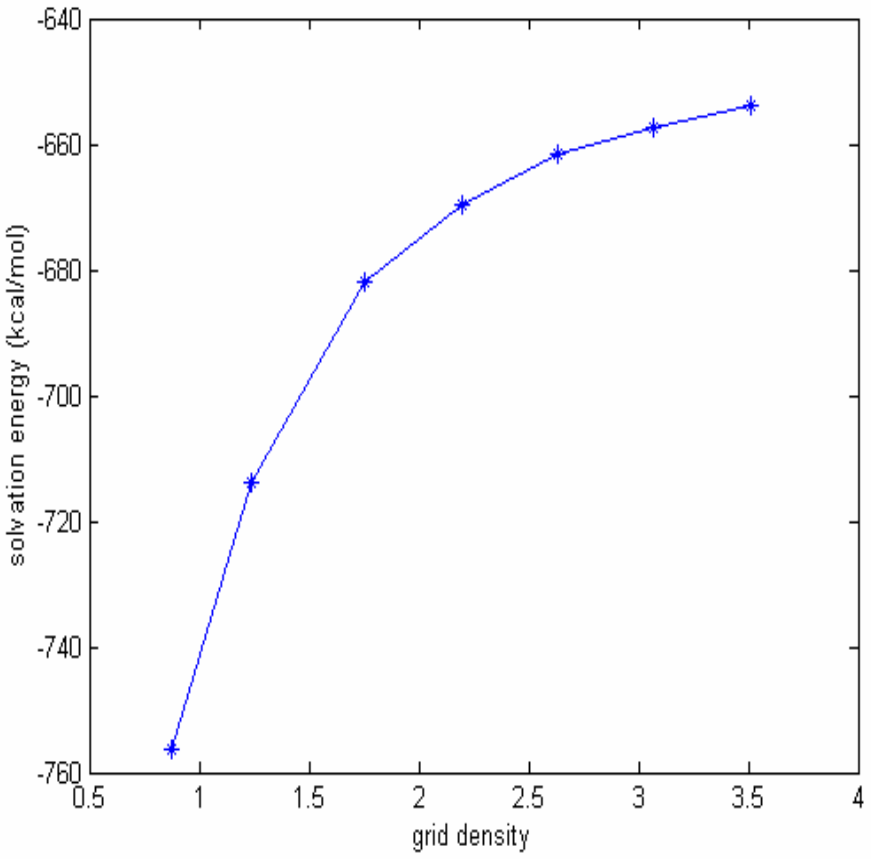
			$E_{solvation}$ (kcal/mol)	
	# of dielectric panels	# of salt panels	pFFT	DelPhi
Water	17204	9330	-3.14	-3.17
TSA	34114	5842	-34.62	-34.75
ECM	82868	18596	-646.42	-653.88

Convergence Results of Ecm Protein

pFFT



DelPhi



Binding energy calculation of a protein-peptide complex

	Energy calculated (kcal/mol)			
	$R_{desolvation}$	$L_{desolvation}$	$(R->L)_{interaction}$	$(L->R)_{interaction}$
pFFT	14.52	24.47	130.80	130.91
DelPhi	14.51	24.47	131.03	131.03

Conclusions and Future work

- Carefully selected Integral Formulation results in Fast Solver for Biomolecule Electrostatics
- Working on coupling to charge optimization problem in drug design
- Extending formulation to include more complicated geometry (inner cavities in macromolecule)
- Fine tuning existing pre-corrected FFT code